

HIGH PRESSURE MAGNETO-OPTICAL STUDIES OF ELECTRONS IN GaSb/AlSb/InAs HETEROSTRUCTURES

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Introduction

Background

☞ GaSb/AlSb/InAs heterostructures \Rightarrow potential for high-speed electronic devices and infrared detectors and sources.

☞ Type-II broken gap (InAs/GaSb) to Type-II staggered (InAs/AlSb) as a function of alloy composition (the crossover occurs at $x = 0.3$ at 1 atm).

☞ Similar tuning of the band alignment can be effected via hydrostatic pressure.

☞ Outstanding problem \Rightarrow *Excess Electrons*.

Possible Sources: Surface donors in the GaSb cap-layer, metastable DX-like bulk states in AlSb, and Tamm-like states at the interfaces with InAs. H. Kroemer, C. Nguyen and B. Brar, J. Vac. Sci. Technol. B **10**, 1769 (1992).

Motivation

To probe the origin of the excess electrons by studying the effects of hydrostatic pressure on the cyclotron resonance (CR) in a single GaSb/AlSb/InAs heterostructure.



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Experimental Setup

FIR magneto-spectroscopy DAC apparatus tunable *in situ*
over 0-150 kbar, 2-300 K, 0-9 T.

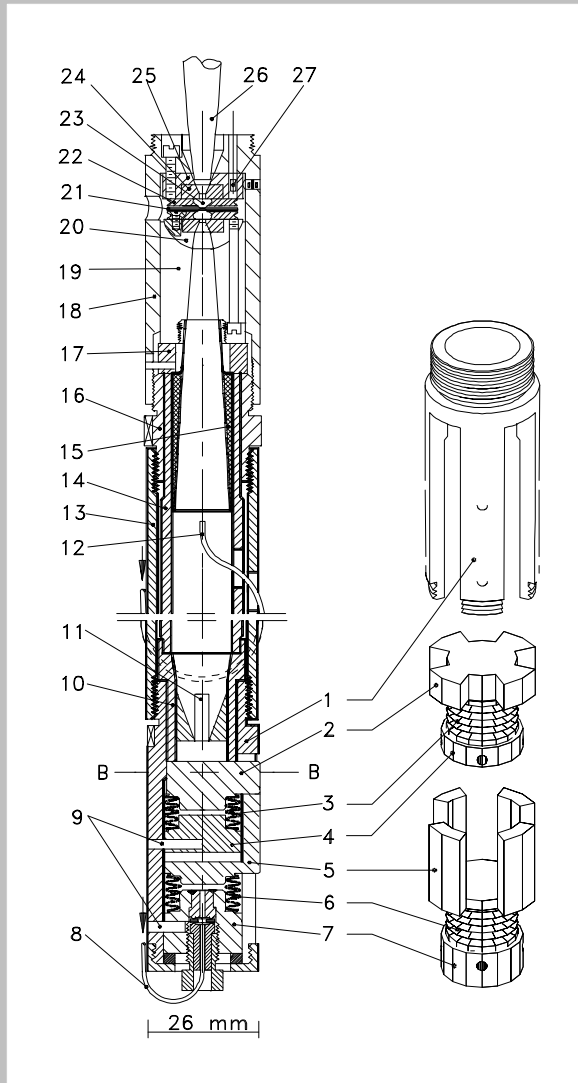
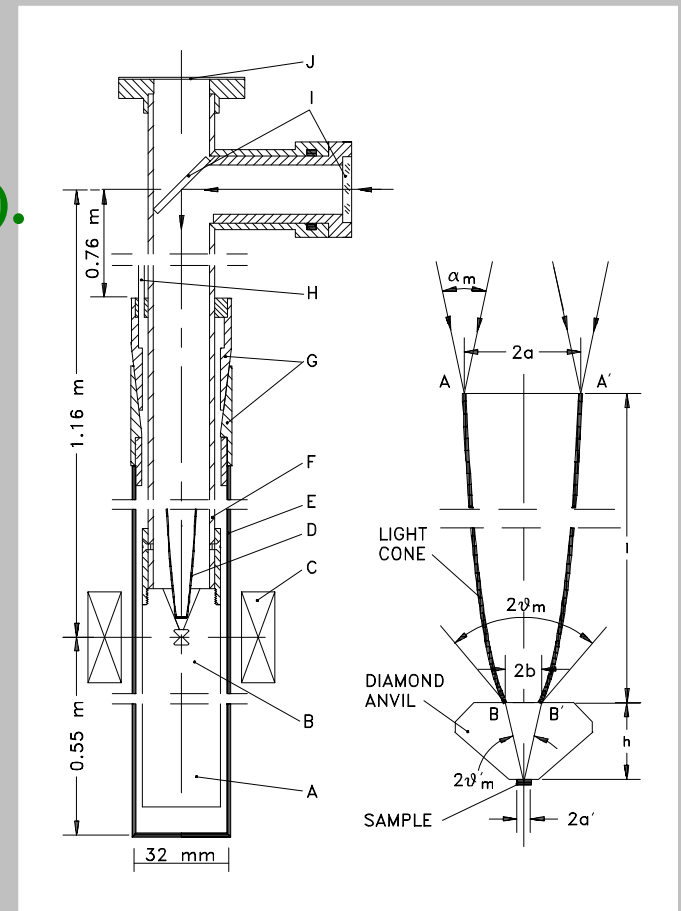
Design Features:

**Multi-bellows
ram(doubles force).**

**Focusing light
cones.**

**Ga:Ge detector
(80-240 cm^{-1}).**

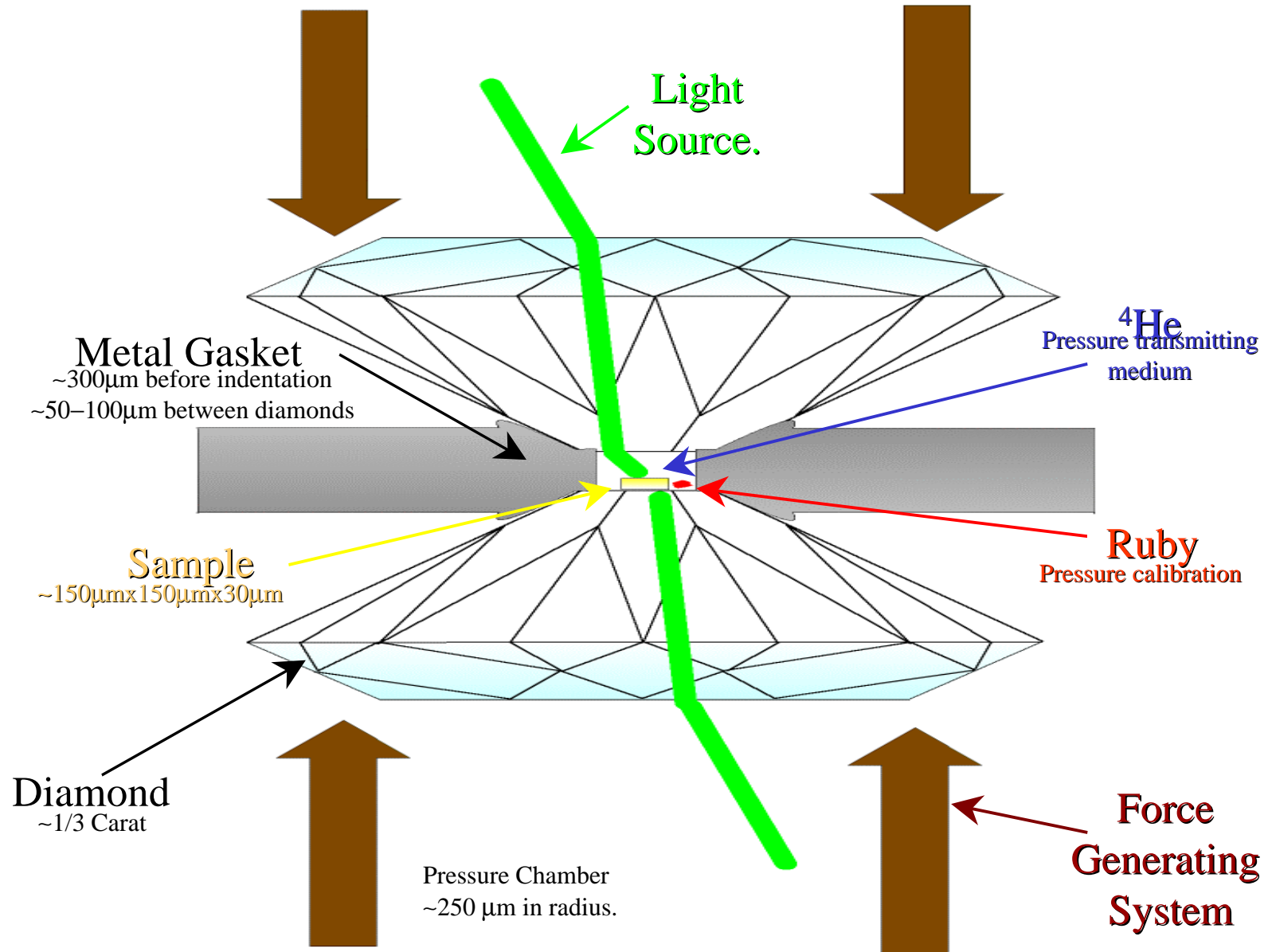
**Visible + FIR
access.**



Multi-bellows ram(US patent 5693345)

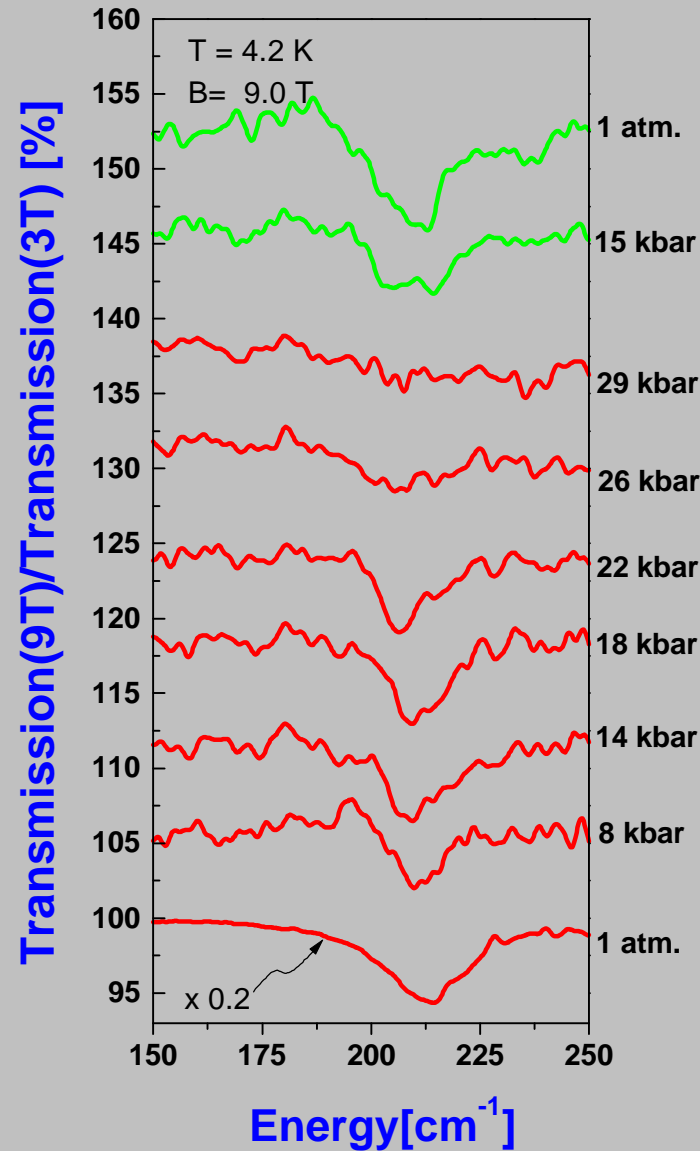
Chen and Weinstein, Rev. Sci. Instrum., **67**, 2883(1996).

Diamond Anvil Cell(DAC)



Schematics of the opposed Bridgman Anvil configuration.

Cyclotron Resonance(CR) Quenching



Effects of pressure on the Fourier transform FIR magneto-transmission at 9T and 4.2 K. The spectra shown in green lines was obtained by releasing the pressure after observing the CR quenching.

Model Solid Theory

☞ **Model-Solid Theory** relates the energy levels in bulk semiconductors to a **common reference level**.

C.G. Van de Walle, Phys. Rev. B **39**, 1871 (1989)

☞ **Strain**, due to applied pressure or to internal mismatch, affects the position of the energy bands.

i =valence band, conduction band

j = AlSb, InAs, GaSb

Average band
energy position.

Deformation
potential.

Fractional
volume change.

$$\Delta E_i^j = a_i^j \times \frac{\Delta \Omega^j}{\Omega^j} = a_i^j \times \left[\frac{\Delta a_x^j}{a_x^j} + \frac{\Delta a_y^j}{a_y^j} + \frac{\Delta a_z^j}{a_z^j} \right]$$

$$\Delta a_x^j = \Delta a_y^j = a_x^j - a_x^{\text{AlSb}}$$

$$\frac{\Delta a_z^j}{a_z^j} = -2 \frac{C_{12}^j}{C_{11}^j} \times \frac{\Delta a_x^j}{a_x^j}$$

☞ **Splitting of the valence band with strain:**

Shear deformation potential for
a strain of tetragonal symmetry.

$$\Delta E_{hh}^j = \frac{\Delta_0^j}{3} - \frac{dE_{001}^j}{2}$$

$J=3/2, m_j=\pm 3/2$

$$\Delta E_{lh}^j = -\frac{\Delta_0^j}{6} + \frac{dE_{001}^j}{4} + \frac{1}{2} \sqrt{(\Delta_0^j)^2 + \Delta_0^j \times dE_{001}^j + \frac{9}{4} (dE_{001}^j)^2}$$

$$dE_{001}^j = 2b^j \times \left(\frac{\Delta a_z^j}{a_z^j} - \frac{\Delta a_x^j}{a_x^j} \right)$$

$J=3/2, m_j=\pm 1/2$

$$\Delta E_{so}^j = -\frac{\Delta_0^j}{6} + \frac{dE_{001}^j}{4} - \frac{1}{2} \sqrt{(\Delta_0^j)^2 + \Delta_0^j \times dE_{001}^j + \frac{9}{4} (dE_{001}^j)^2}$$

$J=1/2, m_j=\pm 1/2$

☞ **Pressure effects**(Murnaghan's equation of state): T.M. Ritter, PhD. Thesis, SUNY at Buffalo (1/97), unpublished.

First pressure derivative
of the bulk modulus.

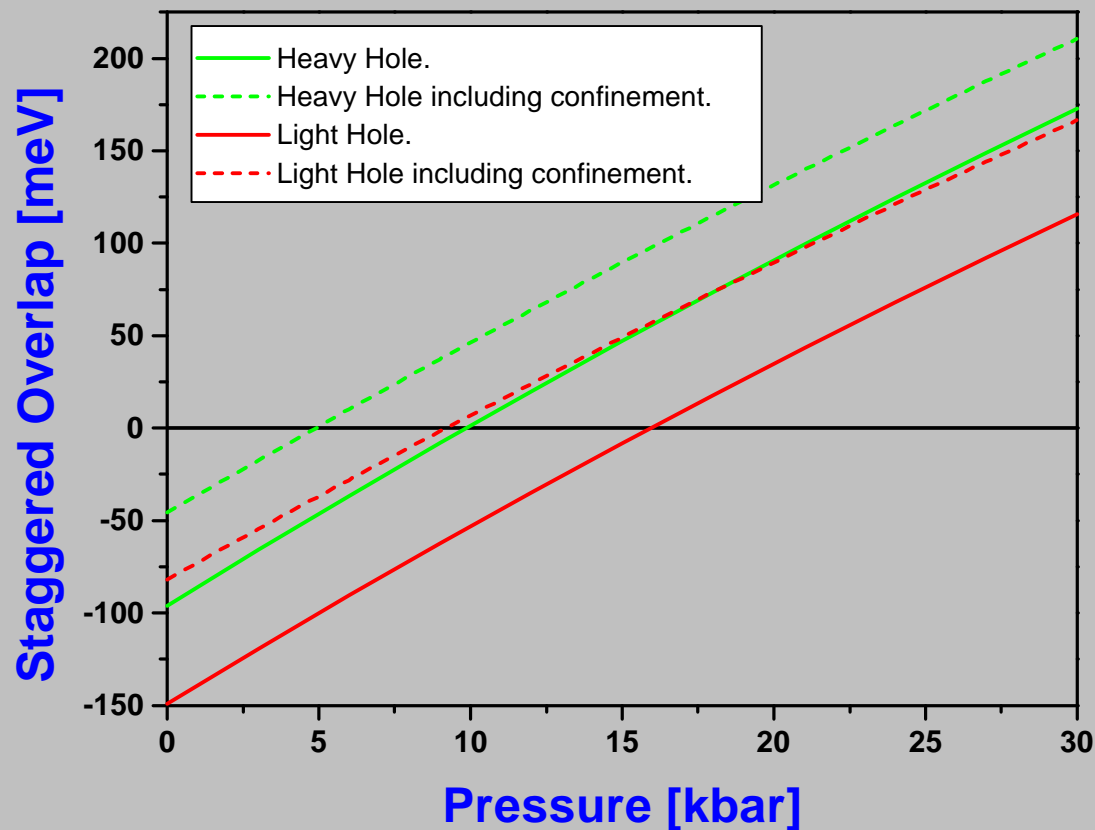
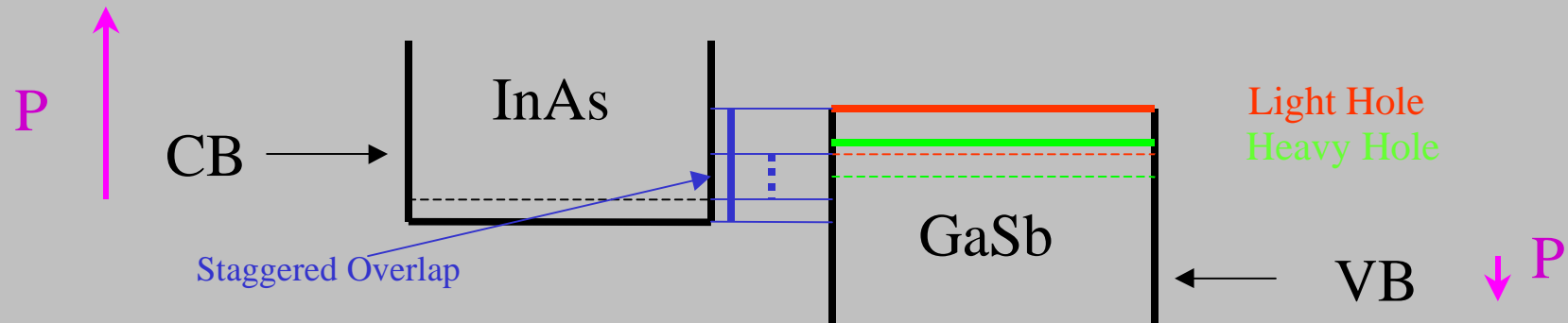
$$a_x^j(p) = a_y^j(p) = a_z^j(p) = a_0^j \times \left[1 + \frac{B_1^j}{B_0^j} p \right]^{\frac{1}{3B_1^j}}$$

1 atm lattice constant

Bulk modulus

$$\Delta a_x^j(p) = \Delta a_y^j(p) = a_x^j(p) - a_0^{\text{AlSb}}$$

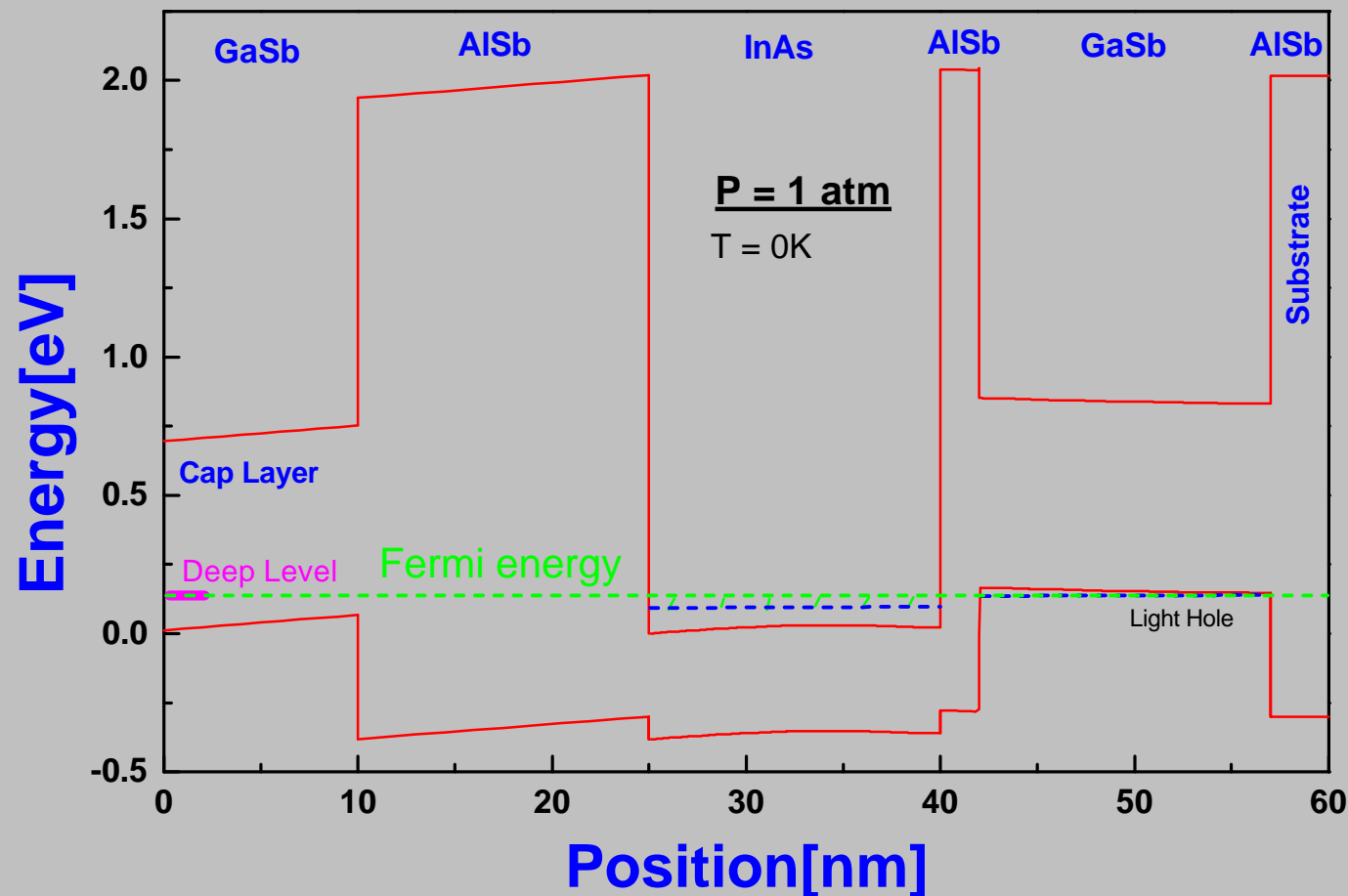
Calculated Results for Flat-Band



Staggered overlap as a function of pressure. The flat band calculation predicts the extinction of the intrinsic electron density at ~9 kbar.

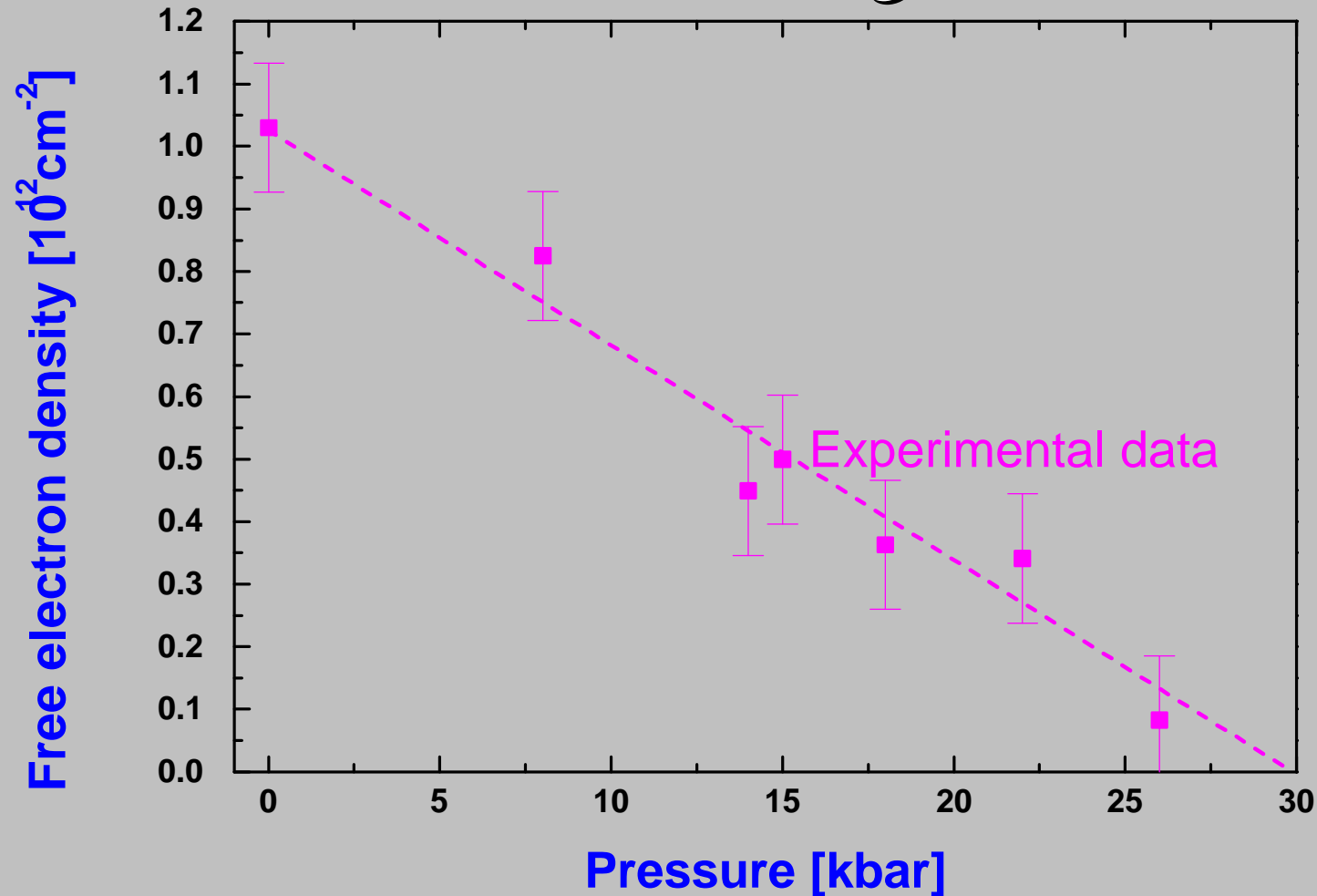
Calculated Results Including Self Consistency

- ➡ **Model Solid Theory.** C.G. Van de Walle, Phys. Rev. B 39, 1871 (1989)
 - ➡ **Variational Solution of Schroedinger's equation for quantum well states.**
 - ➡ **Self-consistent treatment of the electrical potential via Poisson's equation.**
- S. K. Singh, Ph.D. Thesis, SUNY at buffalo, (9/98).



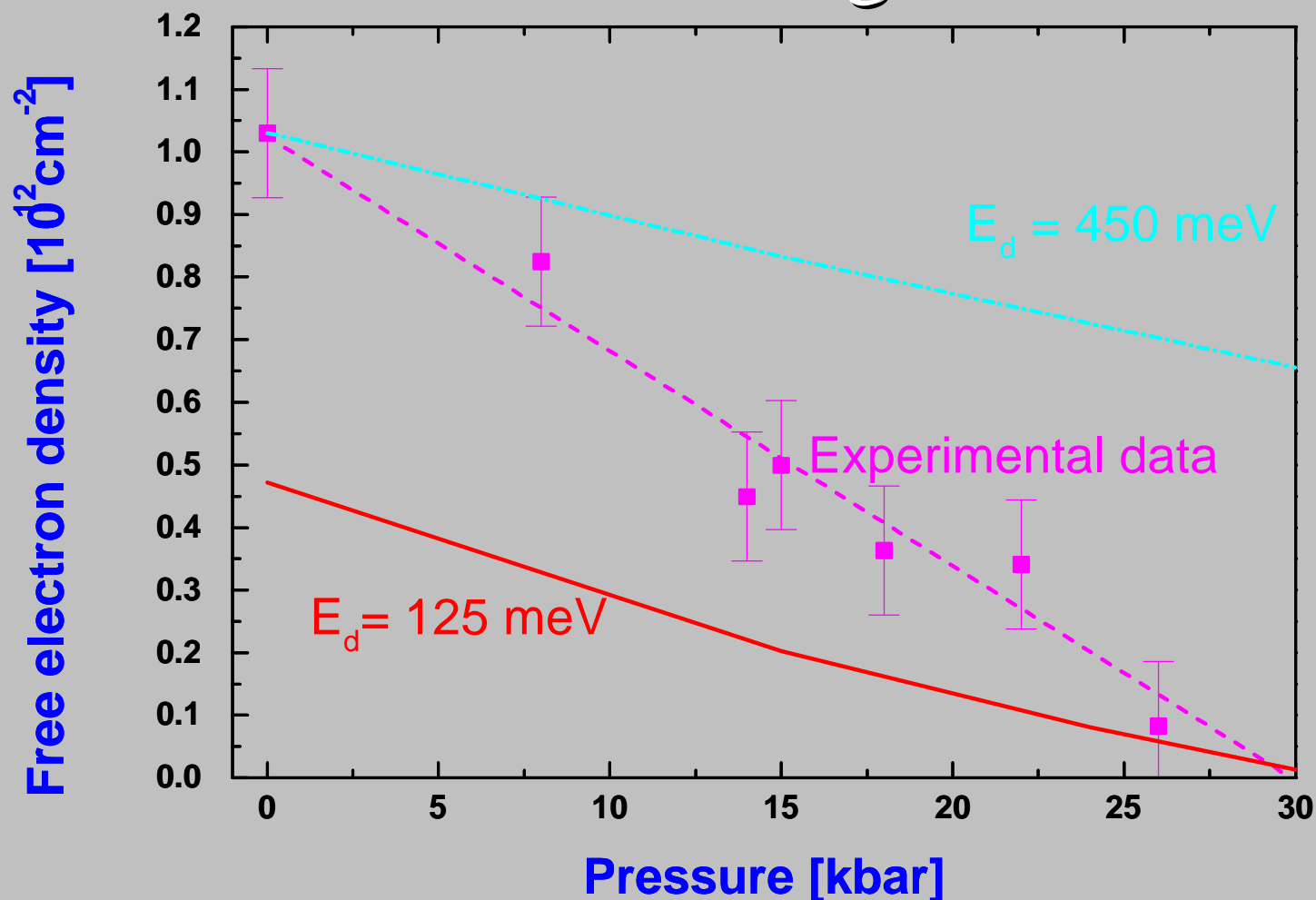
Energy profile calculated at 1atm for the sample studied.

Calculated Results Including Self Consistency



Experimental data calculated from the area of the CR, as a function of pressure, normalized to the 1atm electron density. The 1atm electron density was determined, from the CR strength and spin splitting, to be $(1.0\pm0.1)\times10^{12}\text{cm}^{-2}$.

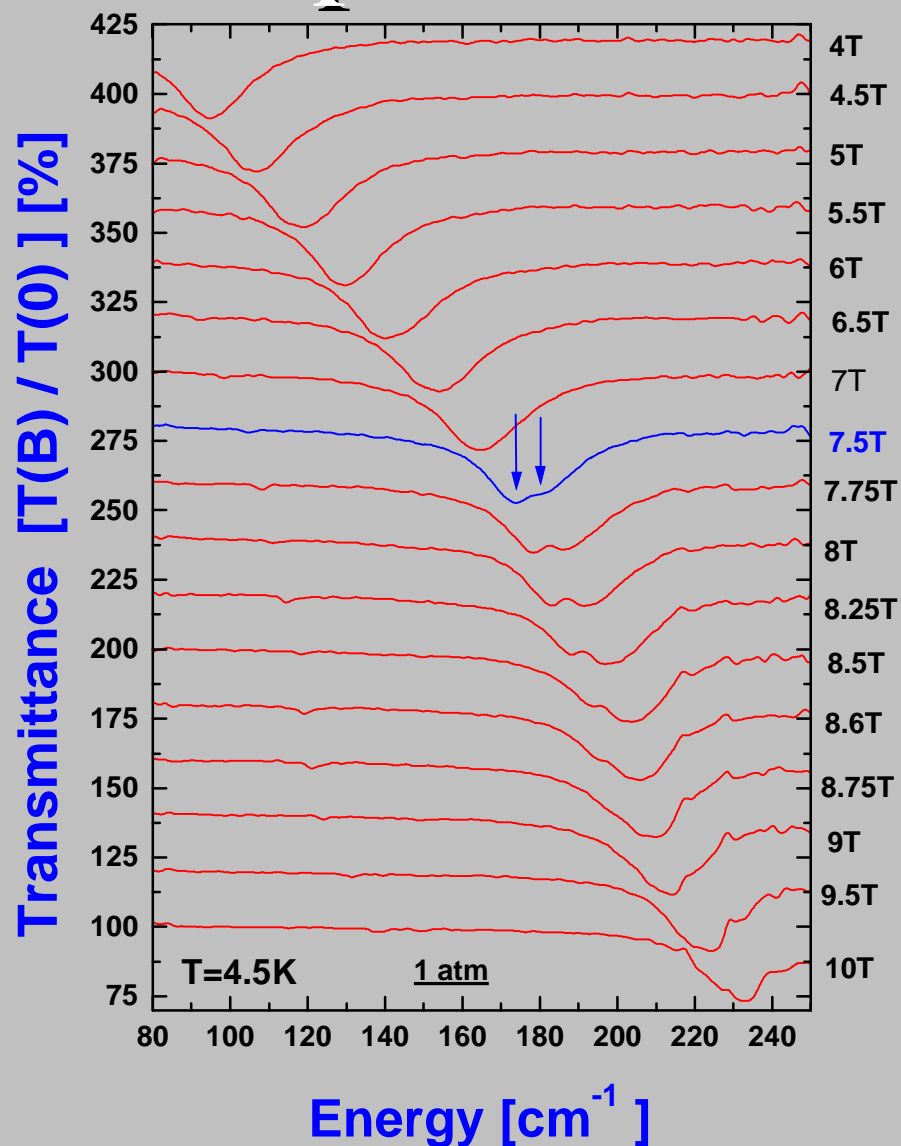
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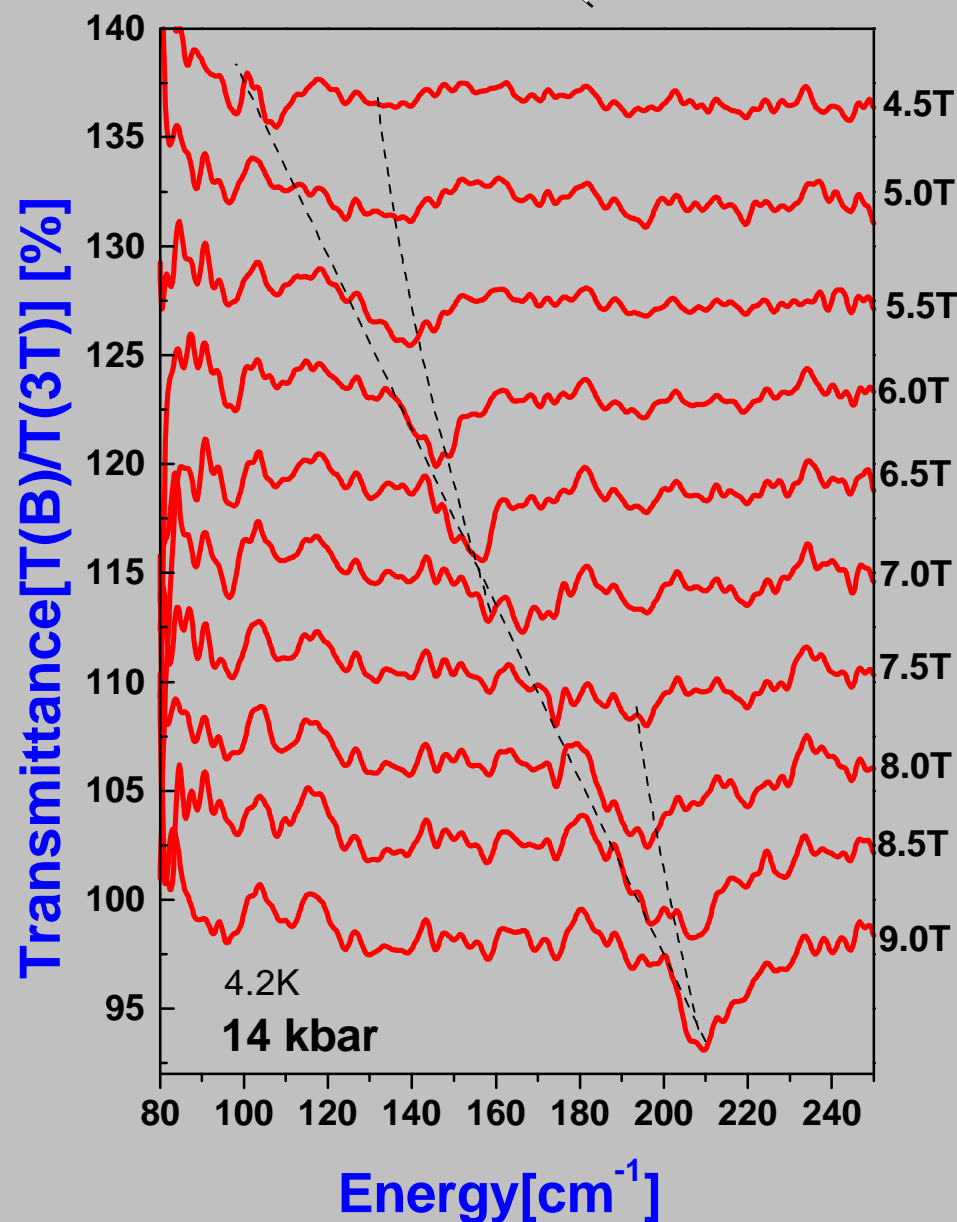
Calculated electron density as a function of pressure assuming a surface donor level at 125 meV and 450 meV respectively from the GaSb valence band.

Atmospheric Data



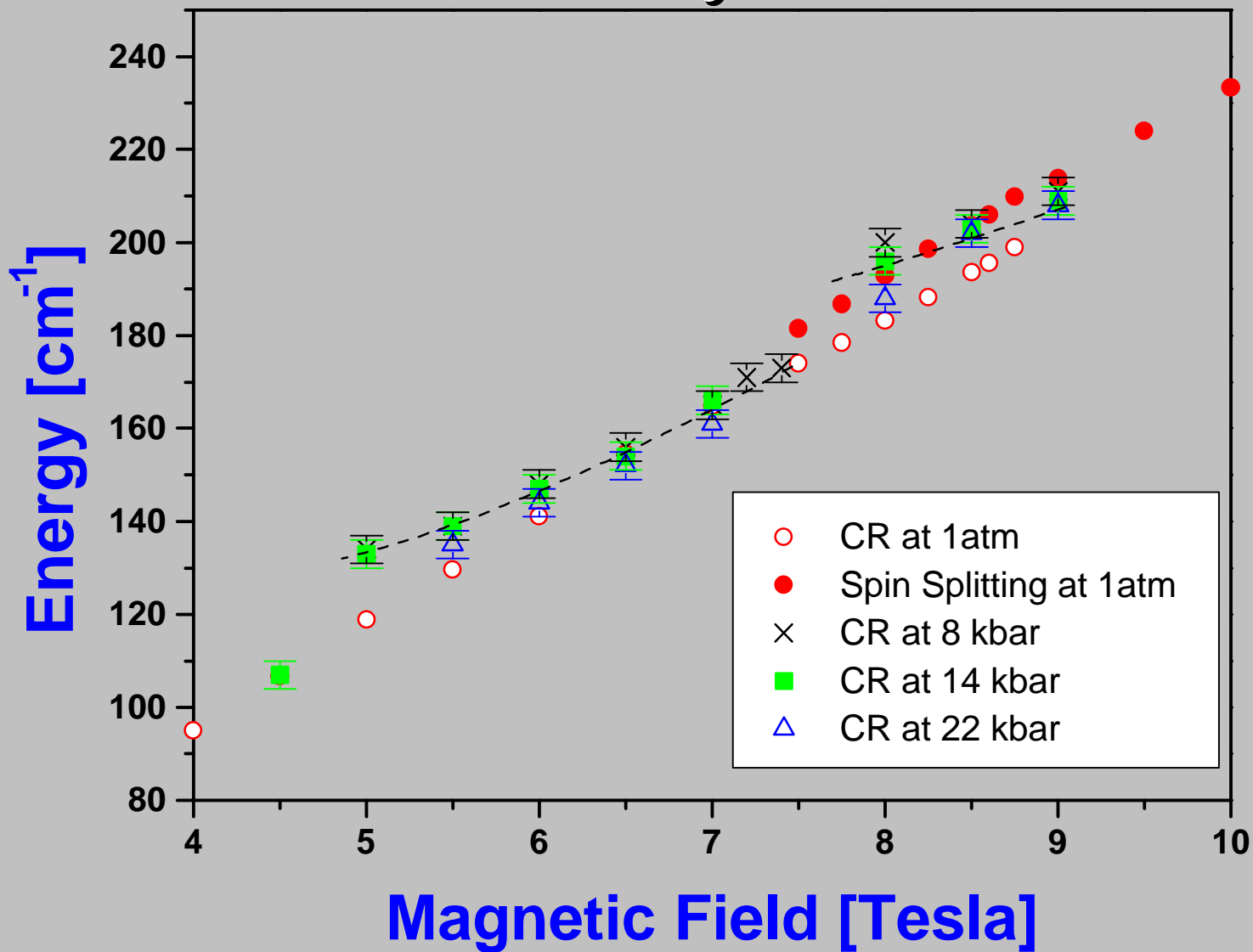
FIR Fourier transform magneto spectroscopy data at 1atm and 4.5K outside the DAC. Notice the spin splitting at ~ 7.5 T shown with the blue trace.

Pressure Data(14 kbar)



FIR Fourier transform magneto spectroscopy data at 14 kbar and 4.2K

Summary Plot



Measured magnetic field dependence of the CR at several different pressures.

Summary and Conclusions

- ☞ CR was recorded as a function of pressure up to 35 kbar, at 4.2 K, and in magnetic fields up to 9T by FIR Fourier transform magneto-transmission spectroscopy in the spectral region 80 --- 250cm⁻¹.
- ☞ CR quenches at (29±3)kbar.
- ☞ Assuming the Fermi level is pinned by surface donor states, we employ a model solid calculation in conjunction with a self consistent treatment.
- ☞ We find that a deep surface state situated 125±25meV above the GaSb valence band can account for the above 29kbar quenching, but can not account for the full electron density at 1 atm.
- ☞ This implies the existence of an additional source/sink of electrons. This could also explain the interesting behavior of the CR with magnetic field if we consider it to be related to metastable states near the InAs/AlSb interface(Tamm-like States).
- ☞ We are pursuing more sophisticated calculations and additional experiments in order to explore this possibility.